clearly shows the differences between the two phases regarding the h.b. distances and angles. In contrast to the $\gamma$ phase the ethylammonium group is disordered in $\beta \mathrm{C} 2 \mathrm{Mn}$ and so is the h.b. scheme. It is the superposition of two orientations. In both of them two h.b.'s run to axial Cl atoms, one to a bridging Cl . This scheme may be designated $a a b$, replacing the term 'monoclinic' (Chapuis, Kind \& Arend, 1976); correspondingly the term 'orthorhombic h.b. scheme' may be replaced by $a b b$; the reason for this new designation is the fact that a 'monoclinic' orientation of the h.b. scheme does not correspond to the symmetry of the structure in cases where $n$ of the general formula is not equal to 1 .

The octahedra layer is puckered and this may be described by specifying the tilt system as $a^{-} a^{-} c^{0}$ (Glazer, 1972; Depmeier, 1977). The tilt angle about the pseudotetragonal $a$ axes is $5^{\circ}$. The four $\mathrm{Cl}(1)$ atoms surrounding the central Mn do not form a regular square but a rectangle with $\mathrm{Cl}(1)-\mathrm{Cl}(1)^{\prime \prime}=3.630$ (3.591 $\AA$ ) and $\mathrm{Cl}(1)^{\prime}-\mathrm{Cl}(1)^{\prime \prime}=3.705$ (3.706 $\AA$ ). The values in brackets refer to the $\gamma$ phase and these values clearly show that the $\gamma \rightarrow \beta$ phase transition mainly affects $b\left[a_{\beta}-a_{v}=0.03, b_{\beta}-b_{v}=0.109 \AA\right.$, cf. also Fig. 6 of Depmeier, Felsche \& Wildermuth (1977)]. This effect should be attributed to the stronger $\mathrm{N}-\mathrm{H}(3) \cdots \mathrm{Cl}(1)$ h.b. in the $\gamma$ phase.* The main

[^0]differences between the $\gamma$ and the $\beta$ phases have been discussed (Depmeier, 1977).

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# 2:3 Complex of Rubidium and 7,7,8,8-Tetracyanoquinodimethane 

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#### Abstract

Rb}_{2}\left(\mathrm{C}_{12} \mathrm{H}_{4} \mathrm{~N}_{4}\right)_{3}\), monoclinic, $a=7.297$ (1), $b$ $=10.392$ (1), $c=21.444$ (2) $\AA, \beta=97.45(1)^{\circ}, Z=2$, $U=1612.3 \AA^{3}, F(000)=772, \mu(\mathrm{Mo})=32.6 \mathrm{~cm}^{-1}$; systematic absences: $h 0 l$ absent for $l$ odd, $0 k 0$ for $k$ odd; space group $P 2_{1} / c$. The final $R_{w}$ is 0.063 for 2750 observed reflexions. Inspection of Weissenberg photographs showed the compound to be isomorphous with $\mathrm{Cs}_{2} \mathrm{TCNQ}_{3}$.


Introduction. As part of a series of crystal structure determinations of TCNQ complexes to study the phase transitions common to most of these salts, the crystal structures of two modifications of RbTCNQ were determined in our laboratory (Hoekstra, Spoelder \& Vos, 1972; van Bodegom, de Boer \& Vos, 1977). The structure determination of a third modification was carried out by Shirotani \& Kobayashi (1973). Recently
we happened to discover crystals of $\mathrm{Rb}_{2} \mathrm{TCNQ}_{3}$ among crystals of RbTCNQ-I.

Cell constants were calculated by a least-squares program from the $\left(\sin ^{2} \theta\right) / \lambda^{2}$ values of 252 reflexions obtained from zero-layer-line Weissenberg photographs superimposed with NaCl reflexion spots $[\lambda(\mathrm{Cu}$ $\left.K \alpha_{2}\right)=1.54434 \AA, \lambda\left(\mathrm{Cu} K \alpha_{1}\right)=1.54050 \AA, a(\mathrm{NaCl})$ $=5.64006 \AA$ at $\left.21^{\circ} \mathrm{C}\right]$. A crystal, $0.2 \times 0.5 \times 0.5$ mm , was mounted on the tip of a glass fibre and aligned along the $b$ axis. Intensities were collected on a threecircle automatic Nonius diffractometer with the $\theta-2 \theta$ method and Zr -filtered Mo radiation. All independent reflexions were measured up to $\theta=30^{\circ}$ at room temperature. Deviations from linearity of the scintillation-counting equipment were kept below $1 \%$ by the use of attenuation filters; every 2 h a reference reflexion was measured three times; corrections were made for intensity changes in the primary beam, for Lorentz and polarization effects and for absorption. The latter correction was calculated with the Busing \& Levy (1957) scheme: $12 \times 10 \times 10$ volume elements and nine accurately located bounding planes were taken into account. Of the 4690 reflexions measured, 3586 had a net intensity greater than zero. A zero

Table 1. Final positional parameters $\left(\times 10^{4}\right)$
Hydrogen atoms were constrained relative to their respective $C$ atoms. For numbering of the atoms see Fig. 1. Numbers in parentheses are the estimated standard deviations in the last significant digits.

|  | $x$ | $y$ | $z$ |
| :--- | ---: | ---: | ---: |
|  |  |  |  |
| $\mathrm{Rb}(1)$ | $5111(1)$ | $2468(1)$ | $2634(0)$ |
| $\mathrm{N}(2)$ | $-3832(8)$ | $125(7)$ | $1712(3)$ |
| $\mathrm{N}(3)$ | $2151(8)$ | $108(7)$ | $2406(3)$ |
| $\mathrm{N}(4)$ | $4606(8)$ | $2996(7)$ | $-962(3)$ |
| $\mathrm{N}(5)$ | $-1168(8)$ | $2964(6)$ | $-1757(3)$ |
| $\mathrm{N}(6)$ | $-4280(8)$ | $3402(7)$ | $1264(3)$ |
| $\mathrm{N}(7)$ | $1545(8)$ | $3528(6)$ | $2065(3)$ |
| $\mathrm{C}(8)$ | $-2339(10)$ | $266(7)$ | $1629(3)$ |
| $\mathrm{C}(9)$ | $-490(8)$ | $453(6)$ | $1501(3)$ |
| $\mathrm{C}(10)$ | $956(10)$ | $265(6)$ | $2013(4)$ |
| $\mathrm{C}(11)$ | $-1447(8)$ | $1114(6)$ | $393(3)$ |
| $\mathrm{C}(12)$ | $-64(8)$ | $916(6)$ | $919(3)$ |
| $\mathrm{C}(13)$ | $1795(8)$ | $1240(6)$ | $835(3)$ |
| $\mathrm{C}(14)$ | $2214(8)$ | $1721(6)$ | $272(3)$ |
| $\mathrm{C}(15)$ | $828(8)$ | $1904(6)$ | $-249(3)$ |
| $\mathrm{C}(16)$ | $-1046(8)$ | $1585(6)$ | $-165(3)$ |
| $\mathrm{C}(17)$ | $3108(9)$ | $2741(7)$ | $-904(3)$ |
| $\mathrm{C}(18)$ | $1261(7)$ | $2406(7)$ | $-823(3)$ |
| $\mathrm{C}(19)$ | $-91(8)$ | $2706(7)$ | $-1340(3)$ |
| $\mathrm{C}(20)$ | $-2804(9)$ | $3668(7)$ | $1209(3)$ |
| $\mathrm{C}(21)$ | $-930(8)$ | $4014(6)$ | $1129(3)$ |
| $\mathrm{C}(22)$ | $433(9)$ | $3760(7)$ | $1658(3)$ |
| $\mathrm{C}(23)$ | $-1853(8)$ | $4679(6)$ | $33(3)$ |
| $\mathrm{C}(24)$ | $-473(8)$ | $4513(6)$ | $575(3)$ |
| $\mathrm{C}(25)$ | $1417(8)$ | $4860(6)$ | $512(3)$ |
| $\mathrm{H}(26)$ | -2864 | 878 | 439 |
| $\mathrm{H}(27)$ | 2882 | 1106 | 1222 |
| $\mathrm{H}(28)$ | 3626 | 1971 | 226 |
| $\mathrm{H}(29)$ | -2130 | 1728 | -552 |
| $\mathrm{H}(30)$ | -3269 | 4420 | 70 |
| $\mathrm{H}(31)$ | 2479 | 4740 | 908 |

weight was given to the reflexions with negative intensity and to 64 high-order reflexions with inaccurate intensities because of inclusion of neighbouring reflexions in the scan.

The coordinates of $\mathrm{Cs}_{2} \mathrm{TCNQ}_{3}$ found by Fritchie \& Arthur (1966) were taken as a starting point for the refinement. The parameters of one Rb atom, one TCNQ group at a general position and one TCNQ around an inversion centre needed to be determined. We refined by blocked-matrix least squares with the XRAY system (1973) set of programs. One block was used for scale and overall temperature factor, one for Rb and for each of the TCNQ groups. Only the 2750 reflexions with $\left|F_{o}\right| \geq 2 \sigma\left(F_{o}\right)$ were taken into account with weights $w=\left[\sigma^{2}\left(F_{o}\right)+\left.E \backslash F_{o}\right|^{2}\right]^{-1}, \sigma\left(F_{o}\right)$ being the standard deviation in $\left|F_{o}\right|$ due to counting statistics and errors in the filter factors; the constant $E=0.00025$ was chosen so as to make $\left\langle w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2}\right\rangle$ as similar as possible for different values of $\left|F_{o}\right|$ and $(\sin \theta) / \lambda$. For the non-hydrogen atoms the scattering curves were taken from Cromer \& Mann (1968) and for H from Stewart, Davidson \& Simpson (1965). For Rb, dispersion corrections as given by Dauben \& Templeton (1955) were applied. H atoms were constrained at $1.08 \AA$ relative to their respective C atoms. For the non-hydrogen atoms anisotropic and for H isotropic parameters were refined. Residual $R_{w}=\left[\sum w\left(\left|F_{o}\right|-\right.\right.$ $\left.\left.\left|F_{c}\right|\right)^{2} / \sum w\left|F_{o}\right|^{2}\right]^{1 / 2}=0.063$, goodness-of-fit $G=$ $\left[\sum w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2} /\left(N_{o}-N_{v}\right)\right]^{1 / 2}=3.31$ with number of observed reflexions $N_{o}=2750$ and number of variables $N_{v}=233$. Final positional parameters are given in Table 1.*

Discussion. The structure is isomorphous with that of $\mathrm{Cs}_{2} \mathrm{TCNQ}_{3}$ (Fritchie \& Arthur, 1966). Both the TCNQ groups and the Rb atoms are arranged in stacks along b. As in $\mathrm{Cs}_{2} \mathrm{TCNQ}_{3}$ there are two types of TCNQ: one group at a centre of symmetry (centrosymmetric c) and another at a general position (non-centrosymmetric nc), the arrangement along $\mathbf{b}$ being nc...c...nc...nc...c...nc. The separation between adjacent TCNQ moieties in the stack, which is a frequently used parameter in the description of structures of TCNQ compounds, is not strictly defined in the present compound as best planes calculated for (parts of) the TCNQ moieties are not parallel. We define the separation between adjacent TCNQ groups as the difference in distance to the origin of the best leastsquares planes through the quinodimethane groups of the molecules considered. For adjacent molecules of mixed type nc...c the calculated distance is $3.26 \AA$

[^1]

Fig. 1. Molecular geometry of the TCNQ units in $\mathrm{Rb}_{2} \mathrm{TCNQ}_{3}$, at $21^{\circ} \mathrm{C}$. Values for bond lengths $(\AA)$ are not corrected for libration effects. Numbering chosen according to Fritchie \& Arthur (1966).
and for adjacent non-centrosymmetric molecules $3 \cdot 12$ $\AA$. The modes of overlap of TCNQ molecules in the stacks are of the same type as in $\mathrm{Cs}_{2} \mathrm{TCNQ}_{3}$. For both

TCNQ groups bond angles and lengths are given in Fig. 1. The Rb ions are surrounded by a distorted cube of eight N atoms with $\mathrm{Rb} \cdots \mathrm{N}$ distances of 3.291 (7), $3 \cdot 263$ (7), $3 \cdot 179$ (6), 2.941 (6), $3 \cdot 118$ (6), $2 \cdot 892$ (6), 3.286 (7) and 3.401 (7) $\AA$.

A determination of the structure at liquid-nitrogen temperature is in progress. An account of the lowtemperature work with full discussion of structure, including structural changes due to differences in atomic radii and temperature, will be published in a following paper.

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[^0]:    * This correlation between h.b. strength and lattice constants seems to be a valuable tool to explain phase transitions in compounds $\left(\mathrm{C}_{n} \mathrm{H}_{2 n+1} \mathrm{NH}_{3}\right)_{2} \mathrm{MnCl}_{4}$ with $n>3$ where single crystals are not available at the moment.

[^1]:    * Lists of structure factors and thermal parameters have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 33332 ( 14 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 13 White Friars, Chester CH1 1NZ, England.

